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Hartmann wavefront measurements at FLASH

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ABSTRACT

We report on online measurements of photon beam parameters during mirror alignment in the soft x-ray spectral region of FLASH, the free-electron laser in Hamburg. A compact Hartmann sensor operating in the wavelength range from 6 to 35nm was used to determine the wavefront quality of individual free-electron laser (FEL) pulses during the alignment procedure as well as aberrations. Beam characterization and alignment of beamline BL3 was performed with $\lambda_{13.5\text{nm}}/116$ accuracy for wavefront rms (w_{rms}). Second moment beam parameters are computed using a spherical reference wavefront generated by a 5 μm pinhole. The Hartmann sensor was used for alignment of the ellipsoidal focusing mirror of beamline BL3, resulting in a reduction of w_{rms} by 33%.

Keywords: Wavefront, Free-electron laser, Hartmann sensor, soft x-ray

1. INTRODUCTION

FLASH, the free-electron Laser in Hamburg is operating in the soft x-ray spectral region based on the self-amplified spontaneous emission (SASE) process in which laser emission is build up from spontaneous undulator radiation. Due to the stochastic nature of this process the beam characteristics relevant for user experiments differ from pulse to pulse, leading to a strong requirement of single pulse photon diagnostic and online characterization of the parameters.^{1,2}

Hartmann and Hartmann-Shack sensors are used for real-time wavefront detection and beam characterization in the near infrared, visible and soft x-ray spectral region. Both, the directional (wavefront) and intensity (beam profile) distribution of the radiation field are provided for single pulses. This enables the evaluation of paraxial beam parameters such as beam width d , divergence θ , beam propagation factor M^2 , Rayleigh length z_R , waist position z_0 and waist diameter d_0 .³ The prediction of the intensity profile in any plane is found by solving the Fresnel-Kirchhoff integral.⁴ In this paper, we present the results from measurements at the FLASH beamline BL3 performed with a compact Hartmann sensor jointly developed by Laser-Laboratorium Göttingen (LLG) and DESY for photon diagnostics, beamline alignment and monitoring of FEL radiation at FLASH^{5,6}. The determination of beam parameters relevant for many user experiments as well as the alignment of the ellipsoidal focusing mirror is described.

2. OPERATION PRINCIPLE

2.1 The Hartmann wavefront sensor

The setup of our Hartmann sensor (see figure 1) is based on the description of Hartmann⁷ and consists of a pinhole array and a 14 bit camera with a charge-coupled device (CCD) chip in a distance $l = 198,251\text{mm}$ behind the pinhole array. The distance represents a compromise between attainable wavefront sensitivity and spatial resolution given by the pitch of the pinhole array. The CCD chip has a size of 8.9mm x 6.7mm with 6.45 μm x 6.45 μm pixel size and is coated with a phosphor (Gd₂O₂S:Tb, grain size 1-2 μm , central emission wavelength 545nm) for converting the soft x-rays into visible light. The Hartmann plate, a pinhole array, made of a 20 μm thick nickel foil with electroformed holes (75 μm diameter and 250 μm pitch) by Precision Eforming divides the incoming beam into an array of smaller beams. The position and intensity of the beams coming from each sub-aperture is measured with the CCD camera.

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http://photon-science.desy.de/facilities/flash/index_eng.html

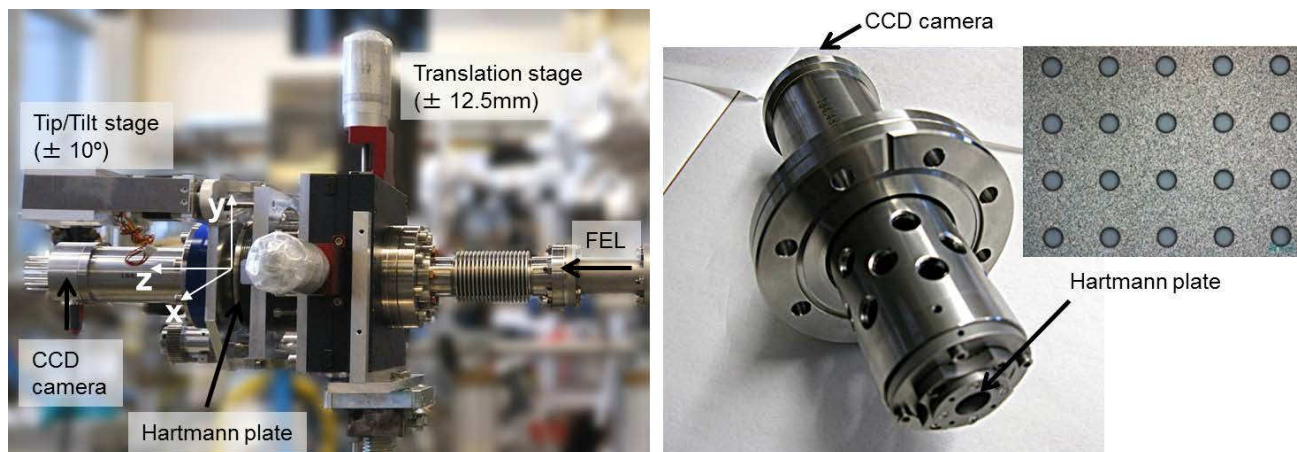


Figure 1. left: Hartmann sensor, right: Inner part of the Hartmann sensor with the Hartmann plate holder. In the inset a light optical microscopic image (magnification 10x) of part of the Hartmann plate with electroformed holes is shown.

The displacement of a spot centroid Δx divided by the distance l yields the local wavefront gradient inside one sub-aperture relative to a known reference wavefront (see figure 2). The wavefront is reconstructed from those local gradients in a modal approach according to the references^{8,9,10}, using 37 Zernike polynomials in the notation of reference¹¹. Summation over all pixel values belonging to the individual sub-apertures samples the intensity distribution or beam profile $I(x, y)$. The Hartmann data, consisting of sampled intensity and wavefront gradients β_x , facilitate to compute the first and second order moments of the spatial (x, y) and angular (u, v) coordinates over the intensity distribution³.

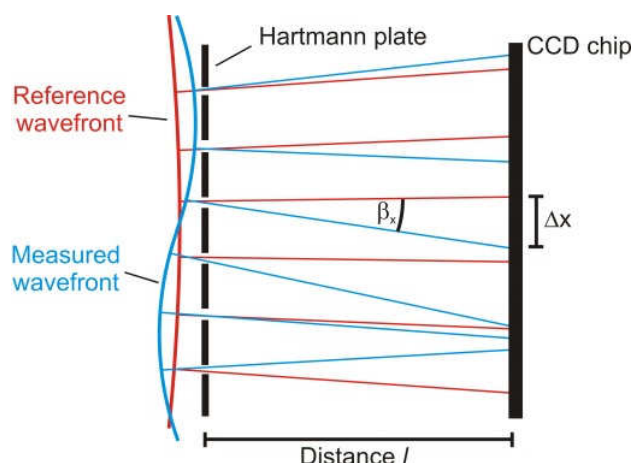


Figure 2. Hartmann principle: The incoming beam (blue) is divided into an array of partial beams by the Hartmann plate. The centroid deviation Δx from the known reference spot position divided by the distance l yields the local wavefront gradient β_x relative to the reference wavefront (red).

This information is sufficient to compute the beam parameters^{3,11} as beam width d , divergence θ , beam propagation factor M^2 , waist diameter d_0 , Rayleigh length z_R and waist position z_0 of paraxial coherent beams. Influences from partial coherence are neglected in this evaluation. The computation of second moment beam parameters from Hartmann data for soft x-ray radiation and the general agreement with caustic scan techniques was reported in reference⁵. Once

intensity and phase of the beam are known from the Hartmann measurements, the Fresnel-Kirchhoff propagation yields the intensity distribution at different positions z .⁴

The area of interest (AOI) for wavefront and beam profile evaluation is selected by clipping noise at a level of 1% of the full dynamic range of the CCD camera. The evaluation radius a is defined by the largest circle inscribed into this AOI. The quality of the wavefront is expressed in terms of wavefront peak-to-valley w_{PV} and wavefront root-mean-square w_{rms} values. Both are calculated after subtraction of tilt and the best-fitting-square.

2.2 Reference wavefront

For wavefront evaluation the Hartmann sensor requires a reference wavefront. Therefore, a $5\mu\text{m}$ pinhole was placed into the FEL beam to generate a spherical wavefront¹². The pinhole was located 110mm upstream of the focal plane at beamline BL2 (ellipsoidal mirror with 2m focal length) and 3592mm in front of the Hartmann plate. FEL pulse train of 30 bunches at a wavelength of 13.5nm were used for calibration measurements. The single pulse energy of about 80 μJ was attenuated using a meshless, 216nm thick Si filter foil to prevent destruction of the pinhole by the FLASH photon beam. For improved signal-to-noise ratio, 100 camera frames, each containing a single pulse train, were averaged to create the reference shown in figure 3. The beam profile was centered on the wavefront sensor with the optical axis perpendicular to the CCD chip. A slightly bigger pinhole (see right bottom in figure 3) in the Hartmann plate defines the sensor orientation to the beam and makes it easier to identify individual sub-apertures.

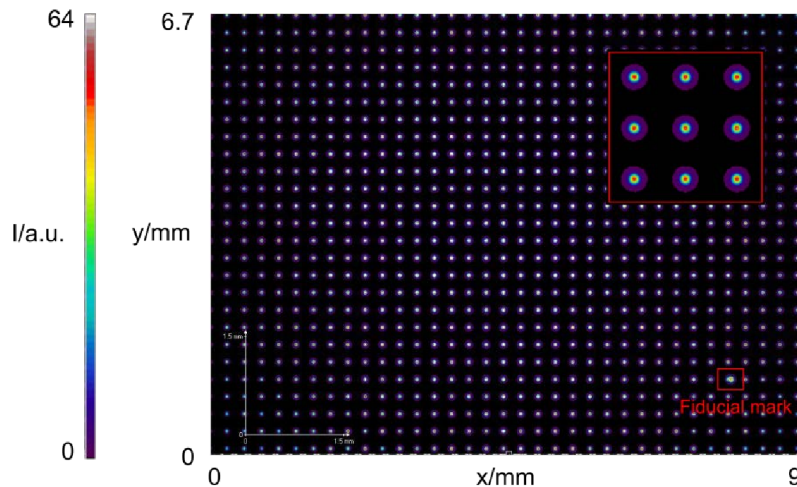


Figure 3. A reference spot pattern is created by diffraction from a $5\mu\text{m}$ pinhole. For alignment purposes, one pinhole (Fiducial mark) bottom right on the Hartmann plate shows a distinct size and shape. It defines the sensor orientation and helps identifying individual apertures.

The average relative single pulse repeatability of the wavefront measurements was determined against the reference shown in figure 3. Therefore a series of 100 single FEL pulses (1 bunch per train) was recorded behind the $5\mu\text{m}$ pinhole at $\lambda=13.5\text{nm}$. Using an evaluation radius of $a=3.07\text{mm}$, a mean deviation for the peak-to-valley value w_{PV} of $\lambda/19$ and for the root-mean-square value w_{rms} of $\lambda/116$ was calculated. Tip/Tilt and defocus are subtracted prior to the computation of the w_{PV} and w_{rms} values.

3. BEAMLINE COMMISSIONING AT BL3

The Hartmann sensor was placed about 850mm downstream of the expected focal plane at FLASH beamline BL3 (see figure 4) to adapt the beam size to the detector area. The measurements were performed in January 2011 using a

wavelength of 13.3nm in single bunch mode with 0.5nC. Two 10mm apertures in an upstream distance of about 48.3m and 52.8m from the focusing ellipsoidal mirror (focal length 2m, 3° grazing incidence angle, amorphous carbon coating) were used. The intensity of about 65μJ was attenuated to about 2μJ per pulse using 0.01mbar Xe in the gas attenuator; the use of transmission filters for further attenuation was not needed. The camera of the wavefront sensor was triggered by the fast shutter which allows us to measure at 10Hz observing online the changes of the Zernike coefficients, w_{rms} , w_{PV} and further beam parameters during the mirror movement.

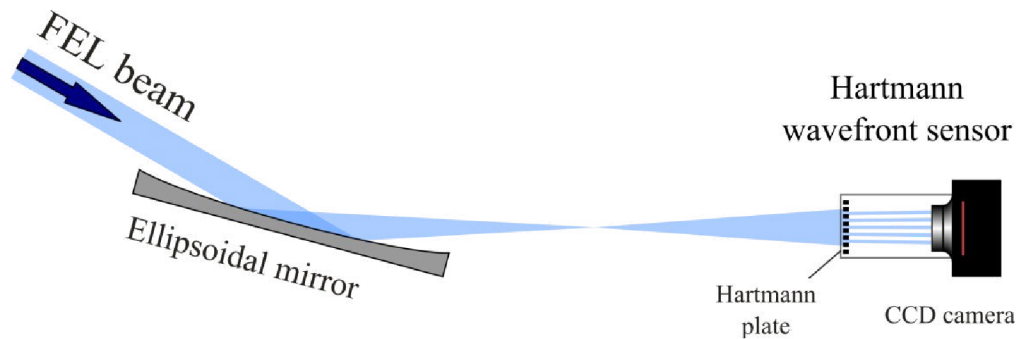


Figure 4. Setup used for the online adjustment of the ellipsoidal grazing incidence focusing mirror of beamline BL3 at FLASH. The ellipsoidal mirror is located in a distance of 72.2m to the source.

The newly installed focusing mirror of beamline BL3 was aligned by successively moving the mirror with respect to the yaw and pitch angles around the pre-aligned position. The goal of alignment procedure with the wavefront sensor is to maximize the Strehl value, which is, according to the Maréchal formula $Strehl \approx 1 - \left(\frac{2\pi}{\lambda} w_{rms}\right)^2$, equivalent to minimizing w_{rms} , with w_{rms} given in nm^{4,13}.

First the yaw angle was aligned such that the astigmatism is oriented along the vertical and horizontal coordinate axis. Therefore, starting at a mirror yaw angle of 0mrad, we scanned in a range between -0.5mrad and 1.5mrad. Afterwards, the pitch angle was adjusted in a range of 0.43mrad and 1.43mrad to minimize the w_{rms} value. For a yaw angle of 0.65mrad we again optimized the pitch angle in a fine-adjustment, finally leading to a reduction of the w_{rms} value by 68% from 4nm to 1.3nm and of w_{PV} by 56% from 27nm to 12nm. In parallel to the mirror movement the astigmatism Zernike coefficients were calculated online to retain information on the success of alignment.

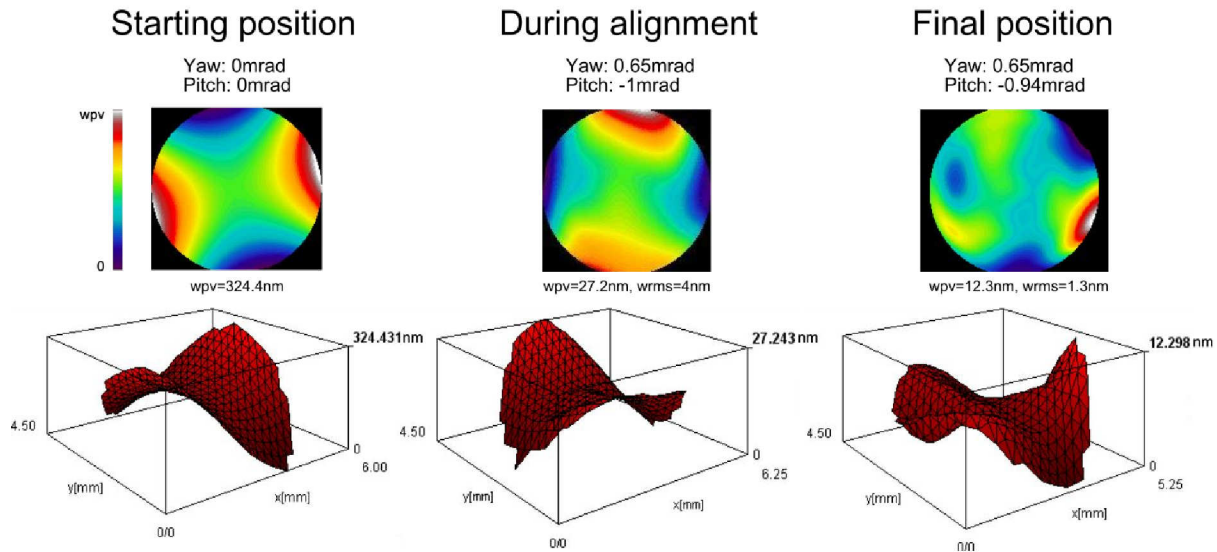


Figure 5. Wavefronts measured at different steps of the alignment procedure of the ellipsoidal focusing mirror at FLASH beamline BL3. The measurements are performed at a wavelength of 13.3nm, in single bunch mode and use of the gas attenuator. The gas attenuator was filled with 0.01mbar Xe to decrease the pulse energy to about 2μJ per pulse.

In figure 5 wavefronts (evaluation radius $a=2.38\text{mm}$) at starting conditions as well as during the alignment procedure and for the final mirror position are shown. The wavefronts shown for the starting conditions and during the alignment are single pulse measurements whereas the wavefront for the final alignment shows an average of 100 single frames on the CCD chip. In table 1 the evaluated beam parameters and astigmatism Zernike aberrations are summarized. The Strehl ratio and the focal spot quality are improved accordingly (see table 1). From the determined final waist position we calculated the focus position to be in a distance of 528.7mm from the flange of the last beamline valve.

Table 1. Beam parameter and Zernike coefficients for different steps of alignment. For the starting position and during the alignment single pulse parameters are given. For the final mirror position parameters are calculated for an average of 100 single pulses.

	Second moment beam parameters at starting position		Second moment beam parameters during alignment		Second moment beam parameters at final position	
	x	y	x	y	x	y
$w_{PV}[\text{nm}]$	324.43		27.24		12.19	
$w_{rms}[\text{nm}]$	-		3.97		1.31	
Beam propagation factor M^2	-		1.94		1.82	
Beam propagation factor M^2_t	-	-	2.25	1.97	1.99	1.83
Beam diameter d [mm]	3.64	3.97	3.86	4.06	3.34	3.85
Waist position z_0 [mm]	-882.85	-813.87	-846.63	-852.61	-848.94	-848.63
Waist diameter d_0 [μm]	107.98	90.37	8.8	7.4	9.1	7.2
Astigmatism $Y^2 - X^2$ [μm]	-0.1338		0.0116		-0.00113	
Astigmatism XY [μm]	0.8632		0.0034		-0.00145	

The Hartmann data recorded before and after the alignment procedure were used to compute beam profiles close to the waist of the FEL beam by numerical propagation via Kirchhoff-Fresnel integral⁵. The corresponding simulated beam profiles are displayed in figure 6. Before the alignment (upper profiles) the strong wavefront astigmatism leads to a

change of the profile orientation before and after the focus, whereas after the alignment a circular beam profile in all planes is observed (lower profiles).

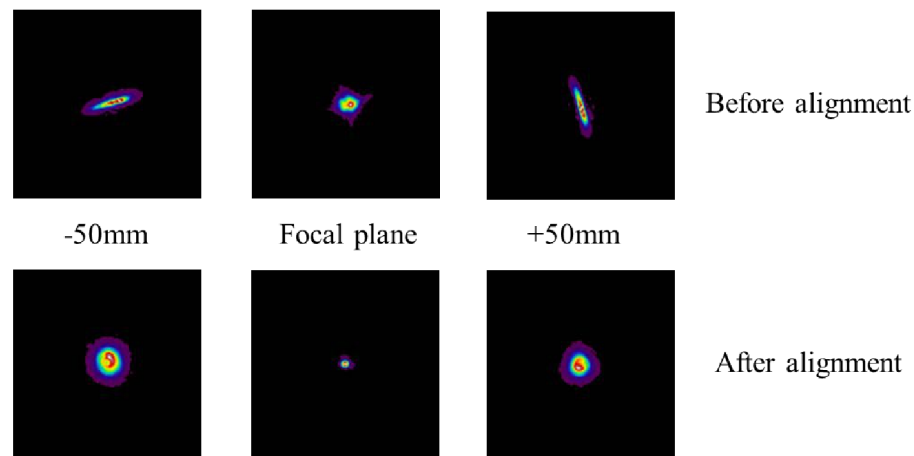


Figure 6. Beam profiles computed for the near-waist region by numerical propagation of Hartmann data before (above) and after (below) the alignment procedure.

4. CONCLUSION

In this paper we report on beam characterization measurements for the free-electron laser FLASH using a Hartmann wavefront sensor for the soft x-ray spectral region with high stability and accuracy. The compact device is mounted on a CF63 vacuum flange and can easily be attached to different beamlines or experimental chambers. This Hartmann sensor with high accuracy is now available for FLASH users and can be applied for evaluation of wavefront, Zernike coefficients, intensity profiles and different beam parameter. A single pulse wavefront repeatability was determined against a spherical reference wavefront, corresponding to $\lambda/116$ at $\lambda=13.5\text{nm}$. We described an alignment procedure of the ellipsoidal focusing mirror of FLASH beamline BL3. The astigmatic distortion dominating the wavefront before the adjustment could strongly ($>60\%$) be reduced to $\lambda/10$ (w_{rms}) at 13.3nm .

ACKNOWLEDGEMENTS

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